

Random measurements are almost maximally incompatible

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Measurement incompatibility

Quantum states and measurements

- Motivation: Classical state $\rightsquigarrow$ **probability distributions**: $p \in \mathbb{R}^d$, $p \geq 0$, $\sum_i p_i = 1$
- Quantum states $\rightsquigarrow$ **density matrices**: $\rho \in M_d(\mathbb{C})$, $\rho \geq 0$, $\text{Tr } \rho = 1$
- Measurement outcomes are labeled $\{1, \dots, k\}$, need to be assigned probabilities
- **Measurements**: Tuples of matrices $(E_1, \dots, E_k)$ such that $(\text{Tr}[E_1\rho], \dots, \text{Tr}[E_k\rho])$ is a probability distribution for all states ρ
 - $\text{Tr}[E_i\rho] \in \mathbb{R} \rightsquigarrow E_i = E_i^*$
 - $\text{Tr}[E_i\rho] \geq 0 \rightsquigarrow E_i \geq 0$
 - $\sum_i \text{Tr}[E_i\rho] = 1 \rightsquigarrow \sum_i E_i = I_d$
- Tuples of PSD matrices summing to identity are called positive operator-valued measures (**POVMs**).

Quantum measurements: Compatibility

- Quantum measurements $\rightsquigarrow$ give the probabilities of the classical outcomes when a quantum state enters a measurement apparatus. Mathematically, measurements are modeled by POVMs.

Definition

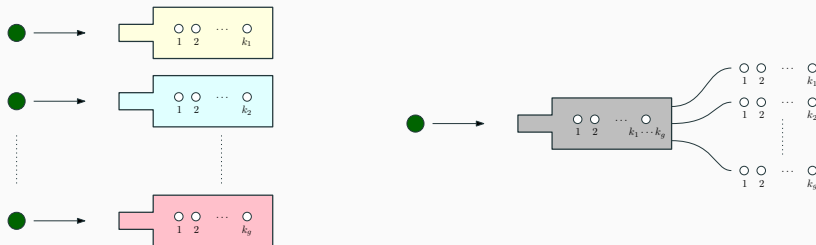
Two POVMs, $E = (E_1, \dots, E_k)$ and $F = (F_1, \dots, F_l)$, are called **compatible** if there exists a third POVM $C = (C_{ij})_{i \in [k], j \in [l]}$ such that

$$\forall i \in [k], \quad E_i = \sum_{j=1}^l C_{ij} \quad \text{and} \quad \forall j \in [l], \quad F_j = \sum_{i=1}^k C_{ij}$$

The definition generalizes to g -tuples of POVMs $E^{(1)}, \dots, E^{(g)}$, having respectively $k_1, \dots, k_g$ outcomes, where the **joint** POVM C has outcome set $[k_1] \times \dots \times [k_g]$.

- Other way to say that: **jointly measurable**

What does it mean?



- Compatible measurements can be simulated by a single joint measurement, by **classically post-processing** its outputs: $E_{i|x} = \sum_{\lambda \in \Lambda} p_{i|\lambda, x} C_{\lambda} \quad \forall i \in [k_x], \forall x \in [g]$.
- Examples:
 1. **Trivial** POVMs $E = (p_i I_d)$ and $F = (q_j I_d)$ are compatible.
 2. **Commuting** POVMs $[E_i, F_j] = 0$ are compatible.
 3. If the POVM E is **projective**, then E and F are compatible iff they commute.

Noisy POVMs

- POVMs can be made compatible by adding **noise**, i.e. mixing in trivial POVMs.
- There are different types of noise, we focus on **white** or **uniform noise**:

Definition

Let $k \in \mathbb{N}$ and let $(E_i)_{i \in [k]}$ be a POVM. Let $t \in [0, 1]$ be a noise parameter. Then, we define the noisy POVM $(E_i^{(t)})_{i \in [k]}$ as

$$E_i^{(t)} := tE_i + (1 - t)\frac{I}{k} \quad \forall i \in [k].$$

Note that $E_i^{(1)} = E_i$ and that $E_i^{(0)} = I/k$.

- Taking $t = 1/2$ suffices to render any pair of dichotomic POVMs compatible.
To see this, choose $C_{ij} := (E_i + F_j)/4$

The compatibility degree

Definition (Compatibility degree of a set of measurements)

Let $d, g \in \mathbb{N}$ and $k_x \in \mathbb{N}$ for all $x \in [g]$. Let $E := \{(E_{i|x})_{i \in [k_x]}\}_{x \in [g]}$ be a set of POVMs. Then, their **compatibility degree** is defined as

$$\tau(E) := \max \left\{ t \in [0, 1] : \left\{ \left(E_{i|x}^{(t)} \right)_{i \in [k_x]} \right\}_{x \in [g]} \text{ is a set of compatible measurements} \right\}.$$

In words: $1 - \tau(E)$ is the minimal amount of noise that makes the measurements in E compatible.

The minimal compatibility degree

Definition (Minimal compatibility degree)

Let $d, g \in \mathbb{N}$ and $k_x \in \mathbb{N}$ for all $x \in [g]$. Then, the **minimal compatibility degree** is defined as

$$\tau(d, g, (k_1, \dots, k_g)) := \min_E \tau(E),$$

where the minimum is taken over all sets E of g POVMs in dimension d , where the x -th POVM has k_x outcomes for all $x \in [g]$.

In words: $1 - \tau(d, g, (k_1, \dots, k_g))$ is the minimal amount of noise that makes any collection of g measurements on a d -dimensional quantum system where the i -th measurement has k_i outcomes compatible.

Some previous results

In [BN18], we have shown that for dichotomic measurements

$$\tau(d, g, (2, \dots, 2)) \geq \frac{1}{\sqrt{g}}$$

and

$$\tau(d, g, (2, \dots, 2)) = \frac{1}{\sqrt{g}} \quad \text{whenever} \quad d \geq 2^{\lceil (g-1)/2 \rceil}.$$

The quantity $\tau(d, g, (2, \dots, 2))$ is the minimal amount of compatibility, so a worst case bound.

Main question

How incompatible are measurements **generically**, i.e., if we look at **random measurements**?

Types of random measurements considered

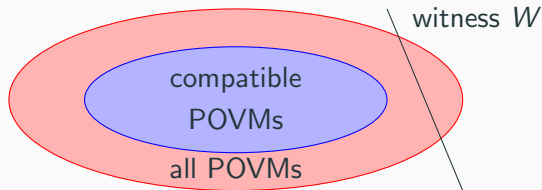
Random POVM model	Number of outcomes	Parameters
Dichotomic PVM	2	rank of the projections
Basis PVM	d	—
Induced POVM	k	dimension of the ancilla space

Types of results:

- **Asymptotic**: Statements which hold with probability 1 in the limit where the underlying dimension d is infinite (tools from free probability)
- **Non-asymptotic**: Statements which hold with probability close to 1 when the underlying dimension d is finite but large (tools from random matrix theory)

Incompatibility witnesses

Incompatibility witnesses



Definition

Let $d, g \in \mathbb{N}$ and $k_x \in \mathbb{N}$ for all $x \in [g]$. An **incompatibility witness** is a g -tuple of operators $W = (W_{i|x})_{i \in [k_x], x \in [g]}$, where $W_{i|x} \in M_d^{\text{sa}}(\mathbb{C})$, such that

$$\max_{E \text{ compatible}} \langle W, E \rangle \leq 1.$$

Here, we have defined

$$\langle W, E \rangle := \sum_{x \in [g]} \sum_{i \in [k_x]} \text{Tr} (W_{i|x} E_{i|x}) .$$

Witnesses for dichotomic measurements

- We can identify a tuple of dichotomic measurements $\{(E_x, I - E_x)\}_{x \in [g]}$ with the set $\{E_x\}$ and set $A_x := 2E_x - I$.
- Then, we can define a **dichotomic incompatibility witness** $W = (W_x)_{x \in [g]}$, where $W_x \in M_d^{\text{sa}}(\mathbb{C})$, as

$$\max_{E \text{ compatible}} \langle W, A \rangle \leq 1.$$

- In [BN22], we have shown that W is an incompatibility witness if and only if there exists a quantum state $\rho \in S_d(\mathbb{C})$ such that

$$\rho - \sum_{x \in [g]} \epsilon_x W_x \geq 0 \quad \forall \epsilon \in \{\pm 1\}^g.$$

Two random dichotomic projective measurements

Constructing random projective measurements

- Fix rank $r \in \mathbb{N}$ and let P_0 be the fixed projection

$$P_0 = \sum_{i=1}^r |i\rangle\langle i| .$$

- Then, take a Haar-random unitary U and define

$$P = UP_0U^* .$$

- Define the measurement to be $(P, I - P)$ and repeat the procedure g times in total with independent unitaries.
- Equivalently, we could have chosen g independent uniformly distributed subspaces of dimension r and used the projections onto them.

Two random projective measurements

Theorem

Fix $0 < \alpha, \beta < 1$ and suppose that

$$\left(\alpha - \frac{1}{2}\right)^2 + \left(\beta - \frac{1}{2}\right)^2 \leq \frac{1}{4}.$$

Let $E, F \subset \mathbb{C}^d$ be independent uniformly distributed subspaces of dimensions $\lfloor \alpha d \rfloor, \lfloor \beta d \rfloor$, respectively. Then, for any $0 < \epsilon \leq 1/\sqrt{2}$,

$$\mathbb{P} \left(\tau(P_E, P_F) \leq \frac{1}{\sqrt{2}} + \epsilon \right) \xrightarrow{d \rightarrow \infty} 1.$$

In this sense, two random projective measurements are **close to being maximally incompatible**.

Proof sketch (1/2)

- Using dichotomic incompatibility witnesses, we can argue that $\tau(P_E, P_F) \leq \tau(V^*P_EV, V^*P_FV)$ for isometries $V : \mathbb{C}^2 \hookrightarrow \mathbb{C}^d$
- We know that $\tau(\sigma_X, \sigma_Z) = \frac{1}{\sqrt{2}}$
- Thus, we are looking for isometries such that

$$V^*(2P_E - I)V = \sigma_Z \quad \text{and} \quad V^*(2P_F - I)V \xrightarrow{d \rightarrow \infty} \sigma_X \text{ in probability.}$$

- We call this the [compression technique](#)

Proof sketch (2/2)

- Let $\gamma = \min(\alpha, \beta, 1 - \alpha, 1 - \beta)$. We can always find orthonormal bases $\{e_1, \dots, e_{\alpha d}\}$ of E and $\{f_1, \dots, f_{\beta d}\}$ of F as well as ordered angles $0 \leq \theta_1 \leq \dots \leq \theta_{\gamma d} \leq \pi/2$ such that, for each $i, j \in [\gamma d]$, $\langle e_i | f_j \rangle = \delta_{i,j} \cos(\theta_i)$. The θ_i 's are called the (non-zero) **principal angles** between the subspaces E and F .
- In other words, P_E and P_F can be block diagonalized such that the blocks are either 1- or 2-dimensional. The 1D blocks have entries 0 or 1, and on the 2D blocks P_E and P_F have the form

$$\begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} \cos(\theta_i)^2 & \cos(\theta_i) \sin(\theta_i) \\ \cos(\theta_i) \sin(\theta_i) & \sin(\theta_i)^2 \end{pmatrix}$$

- This form is well known from Jordan's lemma/Halmos' two projections theorem and [Aub21] proved the distribution of the θ_i for random projections and $d \rightarrow \infty$

Two random projective measurements continued

Theorem

Fix $0 < \alpha < 1$ and suppose that

$$\alpha > \frac{1}{2} \left(1 + \frac{1}{\sqrt{2}} \right) \quad \text{or} \quad \alpha < \frac{1}{2} \left(1 - \frac{1}{\sqrt{2}} \right).$$

Let $E, F \subset \mathbb{C}^d$ be independent uniformly distributed subspaces of dimension $\lfloor \alpha d \rfloor$.

Then, for any $0 < \epsilon \leq 2\sqrt{\lambda_\alpha}\sqrt{1-\lambda_\alpha}/(\sqrt{\lambda_\alpha} + \sqrt{1-\lambda_\alpha})$,

$$\mathbb{P} \left(\tau(P_E, P_F) \leq \frac{1}{\sqrt{\lambda_\alpha} + \sqrt{1-\lambda_\alpha}} + \epsilon \right) \xrightarrow{d \rightarrow \infty} 1, \text{ where } \lambda_\alpha = 4\alpha(1-\alpha),$$

and

$$\mathbb{P} \left(\tau(P_E, P_F) \geq \frac{1}{\sqrt{\lambda_\alpha} + \sqrt{1-\lambda_\alpha}} \right) \xrightarrow{d \rightarrow \infty} 0.$$

For α far from $1/2$, the measurements become **less and less compatible**.

Proof sketch

Recall the [Jordan product criterion](#) from [\[Hei13\]](#):

Lemma

Let $M = (M_i)_{i \in [k]}$ and $N = (N_j)_{j \in [l]}$ be two POVMs on $\mathbb{C}^d$. If, for all $i \in [k]$ and $j \in [l]$,

$$M_i N_j + N_j M_i \geq 0,$$

then M and N are compatible.

We can use this lemma and [\[Aub21\]](#) to show that for $0 \leq t \leq \frac{1}{\sqrt{\lambda_\alpha} + \sqrt{1 - \lambda_\alpha}}$, the POVMs $(P_E^{(t)}, I - P_E^{(t)})$ and $(P_F^{(t)}, I - P_F^{(t)})$ become compatible with probability going to 1 as $d \rightarrow \infty$.

**Many random dichotomic projective
measurements**

Incompatibility of many random dichotomic projective measurements

Theorem

Let $g \leq 3d$ and let $P_1, \dots, P_g$ be random projections of rank $d/2$. Then, with probability at least $1 - e^{-2d}$,

$$\tau(P_1, \dots, P_g) \leq \frac{31}{\sqrt{g}}.$$

- Note that this bound becomes nontrivial only for g large
- Combined with the lower bound $\tau(d, g, (2, \dots, 2)) \geq 1/\sqrt{g}$, this result implies that

$$\tau(d, g, (2, \dots, 2)) = \Theta\left(\frac{1}{\sqrt{g}}\right) \quad \forall d \geq \frac{g}{3},$$

while previously we only knew this to hold for $d \geq 2^{\lceil (g-1)/2 \rceil}$.

Proof sketch

- We need to show that for $\frac{31}{\sqrt{g}} \leq t \leq 1$, the POVMs $(P_x^{(t)}, I - P_x^{(t)})$ are incompatible.
- We need to build an incompatibility witness W such that $\langle W, A \rangle > 1$ for $A_x = 2P_x^{(t)} - I$.
- **Ansatz:** take $W_x = sA_x/d$
- Need to show that for $s \leq c/\sqrt{g}$ and some absolute constant c

$$\frac{I}{d} - \sum_{x=1}^g \epsilon_x W_x \geq 0 \iff \sum_{x=1}^g \epsilon_x A_x \leq \frac{I}{s} \quad \forall \epsilon \in \{\pm 1\}^g,$$

- This can be done using moments of Haar random states and concentration of measure results for sums of sub-Gaussian random variables

Random basis measurements

Random basis measurements

- We consider now the case of measurements in g bases $U_1, \dots, U_g$ of $\mathbb{C}^d$, i.e. whose effects are of the form

$$E_{i|x} = U_x |i\rangle\langle i| U_x^* \quad \forall i \in [d], \forall x \in [g].$$

- Here, U_i are iid Haar random unitary matrices
- Equivalently, we can think of taking random unitaries U_x with columns $\psi_i^{(x)}$ and defining the POVMs to be

$$M_{i|x} = \left| \psi_i^{(x)} \right\rangle \left\langle \psi_i^{(x)} \right| \quad \forall i \in [d]$$

Incompatibility of two random basis measurements

Theorem

Let $U^{(d)}, V^{(d)}$ be sequences of unitary $d \times d$ matrices, such that $U^{(d)}, V^{(d)}$ are independent, and at least one of them is Haar-distributed. Then, for all sequences

$$t_d > \frac{1}{2} \left(1 + \sqrt{\frac{3 \log d}{d}} \right),$$

almost surely as $d \rightarrow \infty$, the measurements

$$\left(t_d U^{(d)} |i\rangle\langle i| U^{(d)*} + (1 - t_d) \frac{I_d}{d} \right)_{i \in [d]} \quad \text{and} \quad \left(t_d V^{(d)} |i\rangle\langle i| V^{(d)*} + (1 - t_d) \frac{I_d}{d} \right)_{i \in [d]}$$

are incompatible.

As for any 2 measurements $\tau(M_1, M_2) \geq 1/2$, it holds that $\tau(B_1, B_2) \xrightarrow{d \rightarrow \infty} \frac{1}{2}$. Hence, two random basis measurements are close to being maximally incompatible.

Incompatibility of many random basis measurements

Theorem

Let $g \geq 2d$ and let $B_1, \dots, B_g$ be random basis measurements. Then, with probability at least $1 - e^{-2g \log(d)}$,

$$\tau(B_1, \dots, B_g) \leq \frac{135 \log(d) - 1}{d - 1}.$$

- This becomes non-trivial only for d large.
- Compatible measurements can be constructed using approximate cloning, thus

$$\tau(B_1, \dots, B_g) \geq \frac{g + d}{g(d + 1)}.$$

- We have shown that up to log factors this bound is tight if g is of order at least d :

$$\tau(d, g, (d, \dots, d)) = \tilde{\Theta} \left(\frac{1}{d} \right) \quad \forall g \geq 2d.$$

Proof sketch

- We want to construct again incompatibility witnesses of the form

$$W_{i|x} = \alpha(U_1, \dots, U_g) U_x |i\rangle\langle i| U_x^*$$

- We can prove that for

$$W_f := \sum_{x \in [g]} W_{f(x)|x},$$

if $\max_{f:[g] \rightarrow [k]} \lambda_{\max}(W_f) \leq \frac{1}{d}$, then W is an incompatibility witness

- Choose $\alpha(U_1, \dots, U_g) = \frac{1}{d \times \eta(U_1, \dots, U_g)}$ with

$$\eta(U_1, \dots, U_g) := \max_{f:[g] \rightarrow [d]} \lambda_{\max} \left(\sum_{x \in [g]} U_x |f(x)\rangle\langle f(x)| U_x^* \right) = \max_{\|\phi\|=1} \sum_{x \in [g]} \|U_x \phi\|_{\infty}^2,$$

- Using facts about Haar-random unitaries and Bernstein inequality with an ε -argument

Random induced POVMs

Constructing random induced POVMs

To obtain a **random POVM** on the system, we proceed as in [HJN20]:

1. Sample a random Haar-distributed isometry $V : \mathbb{C}^d \hookrightarrow \mathbb{C}^k \otimes \mathbb{C}^n$.
2. Define the POVM M on $\mathbb{C}^d$ with k outcomes by

$$M_i := V^*(|i\rangle\langle i| \otimes I_n)V \quad \forall i \in [k].$$

- We construct a **random unital, completely positive map** Φ such that $M_i = \Phi(|i\rangle\langle i|)$
- This construction **induces a probability measures** $\nu_{d,k;n}$ on the set of k -outcome POVMs on a Hilbert space of dimension d .
- The key parameter in this model is the ratio $c := nk/d$, which quantifies the relative size of the ancilla. The value $c = 1$ corresponds to a projective measurement, while $c \rightarrow 0$ corresponds to a trivial POVM $M_i = I/k$ for each $i \in [k]$.
- Although c plays a role analogous to the noise parameter t in the white-noise models discussed previously, there is **no exact correspondence between c and t** .

Incompatibility of random induced POVMs

Theorem

Fix $k, g \in \mathbb{N}$ and consider a sequence of g independent random induced POVMs with k outcomes, distributed along the measure $\nu_{d,k;n_d}$, in the regime where $d \sim ckn_d$ for some fixed parameter c such that

$$c > \frac{4(k-1)g}{(k-1)^2g^2 - 2(k-2)(k-1)g + k^2}.$$

Then, almost surely as $d \rightarrow \infty$, the g POVMs are asymptotically incompatible.

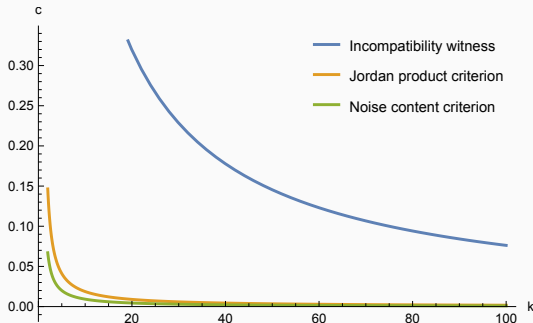
This is proved using incompatibility witnesses and asymptotic freeness of Haar-random unitary matrices and deterministic matrices, implying that $\lim_{d \rightarrow \infty} \lambda_{\max}(W_f)$ can be expressed in terms of a [free convolution](#).

Comparison to compatibility criteria

Noise content criterion [Hei13]: The g POVMs $(M_{i|x})_{i \in [k]}$ are compatible if

$$\sum_{x=1}^g \sum_{i=1}^k \lambda_{\min}(M_{i|x}) \geq g - 1.$$

For $g = 2$, our bound, the noise content and the Jordan product criterion asymptotically have the [same \$1/k\$ scaling](#). However:



Compatibility of random induced POVMs

Theorem

Fix $k, g \in \mathbb{N}$ and consider a sequence of g random induced POVMs with k outcomes, distributed along the measure $\nu_{d,k;n_d}$, in the regime where $d \sim ckn_d$ for some fixed parameter c such that

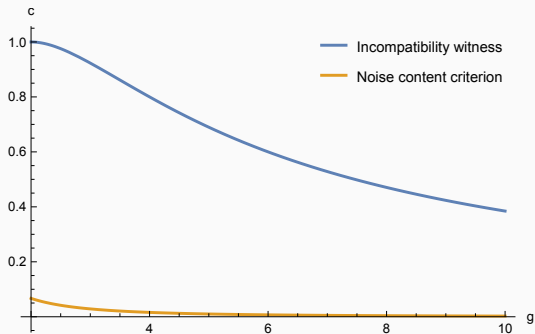
$$c > \frac{1}{g} \times \frac{1}{2g(k-1) + 2\sqrt{(g-1)(k-1)(g(k-1)+1)} - k + 2}.$$

Then, almost surely as $d \rightarrow \infty$, the g POVMs are asymptotically compatible.

This is proved using the noise content criterion and asymptotic freeness.

Comparison to compatibility criterion

For $k = 2$, our upper bound scales as $4/g$, whereas our lower bound scales as $1/(4g^2)$. Numerically:



Open questions

- Can we **extend the compression technique** to $g > 2$ and prove the following conjecture?

Conjecture

Given $g \geq 2$, let $P_1, \dots, P_g$ be g independent Haar-random projections of rank $d/2$ on $\mathbb{C}^d$ and denote by $A_1, \dots, A_g$ their associated dichotomic observables. Then, the following holds in probability:

$$\lim_{d \rightarrow \infty} \tau(A_1, \dots, A_g) = \frac{1}{\sqrt{g}}.$$

- What are **optimal incompatibility witnesses**?
- Other models of randomness? Other forms of non-classicality?

- We have considered random dichotomic projective measurements, random basis measurements, and random induced POVMs
- In high-dimensional regimes and for appropriate parameter choices, [random measurements are typically almost maximally incompatible](#)
- Key techniques: [incompatibility witnesses](#), [compression technique](#)

[[Aub21](#)] G. Aubrun. Principal angles between random subspaces and polynomials in two free projections. *Conflu. Math.*, 13(2), 2021.

[[BN18](#)]: AB and I. Nechita. Joint measurability of quantum effects and the matrix diamond. *J. Math. Phys.*, 59, 2018.

[[BN22](#)]: AB and I. Nechita. A tensor norm approach to quantum compatibility. *J. Math. Phys.*, 63, 2022.

[[Hei13](#)] T. Heinosaari. A simple sufficient condition for the coexistence of quantum effects. *J. Phys. A*, 46(15), 2013.

[[HJN20](#)] T. Heinosaari, M. A. Jivulescu, and I. Nechita. Random positive operator valued measures. *J. Math. Phys.*, 61, 2020.

Our paper:

